

Outlook on the Higgs particles, masses and physical bounds in the Two Higgs-Doublet Model

S.R. Juárez W.* and D. Morales C.†

*Departamento de Física, Escuela Superior de Física y Matemáticas,
Instituto Politécnico Nacional, U.P. “Adolfo López Mateos” C.P. 07738, México D.F.*

P. Kielanowski

*Departamento de Física, Centro de Investigación y Estudios Avanzados
Av. IPN 2508, C.P. 07360, México, D.F.‡*

The Higgs sector of models beyond the standard model requires special attention and study, since through them, a natural explanation can be offered to current questions such as the big differences in the values of the masses of the quarks (hierarchy of masses), the possible generation of flavor changing neutral currents (inspired by the evidence about the oscillations of neutrinos), besides the possibility that some models, with more complicated symmetries than those of the standard model, have a non standard low energy limit. The simplest extension of the standard model known as the two-Higgs-doublet-model (2HDM) involves a second Higgs doublet. The 2HDM predicts the existence of five scalar particles: three neutral (A^0), (h^0 , H^0) and two charged (H^\pm). The purpose of this work is to determine in a natural and easy way the mass eigenstates and masses of these five particles, in terms of the parameters λ_i introduced in the minimal extended Higgs sector potential that preserves the CP symmetry. We discuss several cases of Higgs mixings and the one in which two neutral states are degenerate. As the values of the quartic interactions between the scalar doublets are not theoretically determined, it is of great interest to explore and constrain their values, therefore we analize the stability and triviality bounds using the Lagrange multipliers method and numerically solving the renormalization group equations. Through the former results one can establish the region of validity of the model under several circumstances considered in the literature.

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* Proyecto SIP: 20110368, Secretaría de Investigación y Posgrado, Beca EDI y Comisión de Operación y Fomento de Actividades Académicas (COFAA) del IPN; rebeca@esfm.ipn.mx

† Becario del Consejo Nacional de Ciencia y Tecnología (CONACyT), Mexico; Damoralesc@hotmail.com

‡ Work was partly supported by Polish Ministry of Science and Higher Education Grant N N202 230337; kiel@fis.cinvestav.mx

I. INTRODUCTION

The Standard Model (SM) in high energy physics [1] has been remarkably successful in: describing the properties of elementary particles, predicting the existence of the quarks c , t and b , and the third generation of leptons τ - ν_τ , the existence of the eight gluons, and the weak bosons W^\pm , Z^0 before their discovery, predicting parity violating neutral-weak-currents, and in being consistent with all the experimental results [2, 3]. However, the SM falls short of being a complete theory of the fundamental interactions because of its lack of explanation of the probable unification of the fundamental interactions, the pattern and disparity of the particle masses (mass hierarchy), the origin of the CP violation in nature, the matter-antimatter asymmetry, the pattern of quark mixing, lepton mixing and the reason why there are 3 generations.

As a partial solution to confront these deficiencies, a large number of parameters must be put in “by hand” into the theory (rather than being derived from first principles), such as the three gauge couplings (g_1 , g_2 and g_3), nine fermionic masses (six quarks and three leptons), the Weinberg angle (θ_w), four quark-mixing parameters (CKM) and two more parameters in relation to the Higgs potential (μ and λ).

One of the most subtle aspects of the model is associated with the Higgs sector [4]. The Higgs field and its non-vanishing vacuum expectation value (vev) is the essential ingredient to carry out the spontaneous symmetry breaking (SSB) required to transform the hypothetical massless particles in the Lagrangian into the actual massive physical particles. However, the Higgs particle has not yet been discovered.

In this paper we study the extension of the SM with two Higgs doublets (2HDM) that presents the challenge that the quartic interactions between the scalar doublets are not theoretically determined. This model is studied mainly for three reasons. The first one is that the 2HDM has a much richer Higgs spectrum (3 neutral and 2 charged Higgses) and a different high energy behavior. This makes that a lower mass than in the SM Higgs is permitted. Another reason may be that a different pattern of hierarchy of the Yukawa couplings is possible, because of the presence of two independent vacuum expectation values of the Higgs fields [5]. The third reason is that the Higgs sector of the Minimal Supersymmetric Standard Model (MSSM) contains two Higgs doublets, so the Higgs sectors of the MSSM and the 2HDM are similar and the study of the 2HDM model may give important information on the properties of the Higgs sector in the MSSM [6].

In Section II we introduce the potential for the 2HDM in a special parametrization, and briefly discuss the SSB. In Sections III and IV we present the Higgs mass matrix and its diagonalization method, the mass spectrum, mass eigenstates, and special cases of mixing, where the Higgs masses are simply related to the parameters of the potential. In Section V we obtain and classify the constraints for the quartic couplings derived from the mass formulas, from the vacuum stability principle through the Lagrange multipliers method, and by imposing extreme stability conditions. In Section VI we numerically solve the set of the renormalization group equations from which through triviality principle the physical bounds of the model are determined under different conditions. Finally, Section VII is devoted to the presentation of the results and the conclusions.

II. THE TWO-HIGGS DOUBLET MODEL

In the SM the fermion masses arise, after the SSB, from the couplings between the fermions and a single Higgs doublet. The mass ratio of the b and t quark is of the order of 1/40. To understand in a natural way the origin of this difference in the values of the masses of the third generation of quarks, one can assume the existence of a second Higgs-doublet in the Higgs sector of the SM. In this context one assumes that the quark t obtains its mass through the Φ_1 doublet and the quark b from another doublet Φ_2 [7]. In this way one can explain in a more natural way the hierarchy problem of the Yukawa couplings, as long as the free parameters of the new model acquire the appropriate values.

The Higgs sector of the 2HDM consists of two identical (hypercharge-one) scalar doublets Φ_1 and Φ_2 . There are several proposals for the Higgs potential to describe the physical reality in the framework of the 2HDM [8, 9]. The potential we consider in this paper is compatible with Ref. [10]. It is such that the CP symmetry (charge-conjugation and parity) is conserved, the neutral-Higgs-mediated flavor-changing neutral currents (FCNC) are suppressed in the leptonic sector, and in the quark-sector they are also forbidden by the GIM mechanism [11] in the one loop approximation. In the Lagrangian \mathcal{L} in which we leave out the leptonic terms,

$$\mathcal{L} = \mathcal{L}_{gf} + \mathcal{L}_{Kin} + \mathcal{L}_Y - V \quad (1)$$

the \mathcal{L}_{gf} and \mathcal{L}_{Kin} correspond to kinetic parts of quarks and bosons and they contain the covariant derivatives that provide the interactions among the gauge bosons and the Higgs bosons. They also give rise, after the SSB, to the masses of the gauge bosons (mediators of the electroweak interactions). The fermion masses are generated from the

Yukawa couplings in \mathcal{L}_Y

$$\mathcal{L}_Y = \sum_{i,j} \left(g_{ij}^{(u)} \bar{\psi}_{Li} \Phi_1^c u_{Rj} + g_{ij}^{(d)} \bar{\psi}_{Li} \Phi_2 d_{Rj} \right), \quad (2)$$

between the Higgs bosons and the quarks. In \mathcal{L}_Y , the $g_{ij}^{(u,d)}$ are the Yukawa coupling matrices. The superscripts (u, d) refer to the up and down sectors of quarks, respectively and the subscripts (L, R) correspond to the left handed doublets and right handed singlets in the quark sector. In this paper, we will focus our attention on the potential V .

The Higgs potential

The Higgs potential depends on seven real parameters μ_1^2, μ_2^2 and λ_i ($i = 1, \dots, 5$) from which the five Higgs masses come up after the SSB. The most general renormalizable $SU(2) \times U(1)$ invariant Higgs potential, that preserves a CP and a Z_2 symmetry ($\Phi_1 \rightarrow \Phi_1, \Phi_2 \rightarrow -\Phi_2$) is given by

$$V = \mu_1^2 \Phi_1^\dagger \Phi_1 + \mu_2^2 \Phi_2^\dagger \Phi_2 + \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) \\ + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{1}{2} \lambda_5 \left[(\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2 \right]. \quad (3)$$

For the sake of simplicity a special basis is introduced

$$A = \Phi_1^\dagger \Phi_1, \quad B = \Phi_2^\dagger \Phi_2, \quad C' = D'^\dagger = \Phi_1^\dagger \Phi_2. \quad (4)$$

In this basis

$$V = \mu_1^2 A + \mu_2^2 B + \lambda_1 A^2 + \lambda_2 B^2 + \lambda_3 AB + \lambda_4 C'D' + \frac{1}{2} \lambda_5 [C'^2 + D'^2]. \quad (5)$$

The two Higgs doublets can be represented by eight real fields $\phi_i, i = 1, \dots, 8$,

$$\Phi_1 = \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi_4 \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_5 + i\phi_6 \\ \phi_7 + i\phi_8 \end{pmatrix}. \quad (6)$$

If charge is conserved and there is no CP violation in the Higgs sector, after the SSB, the non-vanishing vacuum expectation values of the fields ϕ_3 and ϕ_7 are real,

$$\langle \phi_3 \rangle = \frac{v_1}{\sqrt{2}}, \quad \langle \phi_7 \rangle = \frac{v_2}{\sqrt{2}}, \quad (7)$$

$$\langle \phi_1 \rangle = \langle \phi_2 \rangle = \langle \phi_4 \rangle = 0, \quad \langle \phi_5 \rangle = \langle \phi_6 \rangle = \langle \phi_8 \rangle = 0. \quad (8)$$

In terms of the fields ϕ_i , the hermitian basis is given by

$$A = \phi_1^2 + \phi_2^2 + \phi_3^2 + \phi_4^2, \quad B = \phi_5^2 + \phi_6^2 + \phi_7^2 + \phi_8^2, \\ C' = \phi_1 \phi_5 + i\phi_1 \phi_6 - i\phi_2 \phi_5 + \phi_2 \phi_6 + \phi_3 \phi_7 + i\phi_3 \phi_8 - i\phi_4 \phi_7 + \phi_4 \phi_8, \\ D' = \phi_1 \phi_5 + i\phi_2 \phi_5 - i\phi_1 \phi_6 + \phi_2 \phi_6 + \phi_3 \phi_7 + i\phi_4 \phi_7 - i\phi_3 \phi_8 + \phi_4 \phi_8. \quad (9)$$

and after the SSB they become

$$\langle A \rangle = \frac{1}{2} v_1^2, \quad \langle B \rangle = \frac{1}{2} v_2^2, \quad \langle C' \rangle = \langle D' \rangle = \frac{1}{2} v_1 v_2. \quad (10)$$

III. THE MASS MATRIX

The conditions for the minimum of the potential are obtained from the vanishing of the first derivatives at the minimum $\frac{\partial V}{\partial \phi_i} \Big|_{\min} = 0$, with the condition that the matrix of the second derivatives at the minimum: $\frac{\partial^2 V}{\partial \phi_i \partial \phi_j} \Big|_{\min}$ is

positive definite. Therefore

$$\begin{aligned} \left. \frac{\partial V}{\partial \phi_i} \right|_{\langle 0 | \phi_i | 0 \rangle} &= \langle 0 | \mu_1^2 \frac{\partial}{\partial \phi_i} A + \mu_2^2 \frac{\partial}{\partial \phi_i} B + 2\lambda_1 A \frac{\partial}{\partial \phi_i} A + 2\lambda_2 B \frac{\partial}{\partial \phi_i} B + \lambda_3 B \frac{\partial A}{\partial \phi_i} + \lambda_3 A \frac{\partial B}{\partial \phi_i} \\ &\quad + \lambda_4 D' \frac{\partial C'}{\partial \phi_i} + \lambda_4 C' \frac{\partial D'}{\partial \phi_i} + \lambda_5 \left[C' \frac{\partial}{\partial \phi_i} C' + D' \frac{\partial}{\partial \phi_i} D' \right] | 0 \rangle = 0 \end{aligned} \quad (11)$$

from which after some simplifications two non trivial equations are obtained

$$\mu_1^2 + \lambda_1 v_1^2 + 2\lambda_T v_2^2 = 0 \text{ or } v_1 = 0, \quad \mu_2^2 + \lambda_2 v_2^2 + 2\lambda_T v_1^2 = 0 \text{ or } v_2 = 0, \quad (12)$$

where

$$\lambda_T \equiv (\lambda_3 + \lambda_4 + \lambda_5). \quad (13)$$

The mass matrix elements are obtained from the equation

$$M_{ij}^2 = \frac{1}{2} \left. \frac{\partial^2 V}{\partial \phi_i \partial \phi_j} \right|_{\phi_3 = \frac{v_1}{\sqrt{2}}, \phi_7 = \frac{v_2}{\sqrt{2}}}, \quad (14)$$

and the explicit form of the matrix of the second derivatives reads

$$\begin{aligned} \frac{\partial^2 V}{\partial \phi_j \partial \phi_i} &= (\mu_1^2 + 2\lambda_1 A + \lambda_3 B) \frac{\partial^2 A}{\partial \phi_j \partial \phi_i} + (\mu_2^2 + 2\lambda_2 B + \lambda_3 A) \frac{\partial^2 B}{\partial \phi_j \partial \phi_i} + 2\lambda_1 \frac{\partial A}{\partial \phi_j} \frac{\partial A}{\partial \phi_i} + 2\lambda_2 \frac{\partial B}{\partial \phi_j} \frac{\partial B}{\partial \phi_i} \\ &\quad + \lambda_3 \left(\frac{\partial A}{\partial \phi_j} \frac{\partial B}{\partial \phi_i} + \frac{\partial B}{\partial \phi_j} \frac{\partial A}{\partial \phi_i} \right) + \lambda_5 \left(\frac{\partial C'}{\partial \phi_j} \frac{\partial C'}{\partial \phi_i} + \frac{\partial D'}{\partial \phi_j} \frac{\partial D'}{\partial \phi_i} \right) + \lambda_4 \left(\frac{\partial C'}{\partial \phi_j} \frac{\partial D'}{\partial \phi_i} + \frac{\partial D'}{\partial \phi_j} \frac{\partial C'}{\partial \phi_i} \right) \\ &\quad + (\lambda_4 D' + \lambda_5 C') \frac{\partial^2 C'}{\partial \phi_j \partial \phi_i} + (\lambda_4 C' + \lambda_5 D') \frac{\partial^2 D'}{\partial \phi_j \partial \phi_i}. \end{aligned} \quad (15)$$

Using Eqs. (12) the 16 non vanishing matrix elements are

$$\begin{aligned} M_{11}^2 &= M_{22}^2 = -\frac{1}{2} (\lambda_4 + \lambda_5) v_2^2, & M_{33}^2 &= 2\lambda_1 v_1^2, & M_{44}^2 &= -\lambda_5 v_2^2, \\ M_{55}^2 &= M_{66}^2 = -\frac{1}{2} (\lambda_4 + \lambda_5) v_1^2, & M_{77}^2 &= 2\lambda_2 v_2^2, & M_{88}^2 &= -\lambda_5 v_1^2, \\ M_{15}^2 &= M_{51}^2 = M_{26}^2 = M_{62}^2 = \frac{1}{2} (\lambda_4 + \lambda_5) v_1 v_2, \\ M_{37}^2 &= M_{73}^2 = (\lambda_3 + \lambda_4 + \lambda_5) v_1 v_2, & M_{48}^2 &= M_{84}^2 = \lambda_5 v_1 v_2. \end{aligned} \quad (16)$$

Diagonalization of the mass matrix

The Higgs masses and the Higgs mass-eigenstates are obtained after a suitable diagonalization of the matrix in Eq. (14). The diagonalization of the matrix whose elements are given in Eq. (16) is performed in two steps.

A block ordering is performed and a diagonalization of each block is carried out. The blocks are obtained by means of the application of two consecutive unitary transformations $U_1 = U_1^\dagger$ and $U_2 = U_2^\dagger$

$$(M_{ij}^2)_B = U_2 U_1 M_{ij}^2 U_1^\dagger U_2^\dagger. \quad (17)$$

The non vanishing matrix elements of the unitary transformations are

$$\begin{aligned} (U_1)_{11} &= (U_1)_{25} = (U_1)_{33} = (U_1)_{44} = (U_1)_{52} = (U_1)_{66} = (U_1)_{77} = (U_1)_{88} = 1, \\ (U_2)_{11} &= (U_2)_{22} = (U_2)_{33} = (U_2)_{47} = (U_2)_{55} = (U_2)_{66} = (U_2)_{74} = (U_2)_{88} = 1. \end{aligned}$$

After carrying out both transformations, the 8×8 matrix becomes

$$(M_{ij}^2)_B = \begin{pmatrix} M_{11}^2 & M_{15}^2 & 0 & 0 & 0 & 0 & 0 & 0 \\ M_{51}^2 & M_{55}^2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & M_{33}^2 & M_{37}^2 & 0 & 0 & 0 & 0 \\ 0 & 0 & M_{73}^2 & M_{77}^2 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & M_{22}^2 & M_{26}^2 & 0 & 0 \\ 0 & 0 & 0 & 0 & M_{62}^2 & M_{66}^2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & M_{44}^2 & M_{48}^2 \\ 0 & 0 & 0 & 0 & 0 & 0 & M_{84}^2 & M_{88}^2 \end{pmatrix}. \quad (18)$$

The matrix in Eq. (18), is ready to easily perform the total diagonalization

$$(M_{ij}^2)_B = \begin{pmatrix} M_{11}^2 & M_{15}^2 \\ M_{51}^2 & M_{55}^2 \end{pmatrix} \oplus \begin{pmatrix} M_{33}^2 & M_{37}^2 \\ M_{73}^2 & M_{77}^2 \end{pmatrix} \oplus \begin{pmatrix} M_{22}^2 & M_{26}^2 \\ M_{62}^2 & M_{66}^2 \end{pmatrix} \oplus \begin{pmatrix} M_{44}^2 & M_{48}^2 \\ M_{84}^2 & M_{88}^2 \end{pmatrix}. \quad (19)$$

The next step is to perform the diagonalization of each of the submatrices in Eq. (19).

IV. HIGGS MASS-EIGENSTATES BASIS

Let us now proceed to relate the gauge states with the mass eigenstates.

The scalar fields in Eq. (6) can be represented as

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \phi_1^0 \end{pmatrix}; \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \phi_2^0 \end{pmatrix} \quad (20)$$

where

$$\phi_1^+ = \phi_1 + i\phi_2, \quad \phi_2^+ = \phi_5 + i\phi_6, \quad \phi_1^0 = \phi_3 + i\phi_4, \quad \phi_2^0 = \phi_7 + i\phi_8, \quad (21)$$

and

$$\phi_3 = \frac{v_1}{\sqrt{2}} + h_1, \quad \phi_4 = \eta_1, \quad \phi_7 = \frac{v_2}{\sqrt{2}} + h_2, \quad \phi_8 = \eta_2. \quad (22)$$

Now, the physical fields (mass eigenstates) and the Goldstones (massless eigenstates) are obtained from the gauge eigenstates by a unitary transformation that diagonalizes the corresponding submatrices Eq. (19), in the following way

$$\begin{pmatrix} H^0 \\ h^0 \end{pmatrix} = U_\alpha \begin{pmatrix} h_1 \\ h_2 \end{pmatrix}, \quad \begin{pmatrix} G^+ \\ H^+ \end{pmatrix} = U_\beta \begin{pmatrix} \phi_1^+ \\ \phi_2^+ \end{pmatrix}, \quad \begin{pmatrix} G^0 \\ A^0 \end{pmatrix} = U_\gamma \begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix}, \quad (23)$$

where

$$U_\alpha = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix}, \quad U_\alpha^\dagger U_\alpha = I \quad (24)$$

and $U_\beta = U_\gamma$ have the same form as U_α . α is the mixing angle between the neutral states ϕ_1^0 and ϕ_2^0 , i.e., ϕ_3 and ϕ_7 , the β angle is the one between the charged states, and γ is related with the CP-Odd states ϕ_4 and ϕ_8 .

$$\tan \alpha = \frac{y}{\sqrt{1+y^2}}, \quad y = \tan 2\alpha = \frac{(\lambda_3 + \lambda_4 + \lambda_5)v_1 v_2}{(\lambda_1 v_1^2 - \lambda_2 v_2^2)}, \quad -\frac{\pi}{2} < \alpha < \frac{\pi}{2}, \quad (25)$$

and

$$\tan \beta = \frac{v_2}{v_1}, \quad v^2 = (v_1^2 + v_2^2), \quad 0 < \beta < \frac{\pi}{2}, \quad \gamma = \beta. \quad (26)$$

The resulting physical particles in the Higgs-sector are: two CP-even-neutral Higgs scalars (H^0, h^0), one CP-odd neutral Higgs scalar (A^0), two charged Higgs bosons (G^\pm), and three Goldstone-bosons (G^\pm, G^0) that contribute to the mass generation of the gauge vector bosons W^\pm and Z^0 , respectively.

The mass formulas

After the complete diagonalization, we obtain the following relations:

1. The mass eigenvalues for (H^0, h^0) are

$$M_{H^0, h^0}^2 = \lambda_1 v_1^2 + \lambda_2 v_2^2 \pm \sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + (v_1 v_2 \lambda_T)^2} > 0, \quad (27)$$

2. The eigenvalues for the mass eigenstates H^\pm and G^\pm are

$$M_{G^\pm}^2 = 0, \quad M_{H^\pm}^2 = -\frac{1}{2} (\lambda_4 + \lambda_5) v^2 > 0, \quad (28)$$

3. Finally, the mass eigenvalues for G^0 and A^0 are

$$M_{G^0}^2 = 0, \quad M_{A^0}^2 = -\lambda_5 v^2 > 0, \quad (29)$$

As expected, after the electroweak symmetry breaking (EWSB) the eight components of the two complex isodoublet fields are transformed into: two charged Higgs bosons H^\pm , three neutral Higgs bosons H^0, h^0, A^0 , and three massless Goldstone fields G^0, G^\pm (which are transformed into the longitudinal components of the gauge bosons W^\pm and Z^0). At this level, the values of M_{A^0} and M_{H^\pm} are not related to the parameters λ_1, λ_2 and λ_3 . This means that, apparently, there is a complete independence between the A^0, H^\pm and the h^0, H^0 , which is not all true.

As in the Standard Model, the values of the quartic couplings are not fixed by the model. To proceed as in the SM [12], to determine the Higgs masses, one has to consider two important physical principles. The vacuum stability constrains the values for the quartic couplings. To have a complete view, we invert the former equations to express the quartic parameters in terms of the masses of the Higgs fields.

$$\lambda_1 = \frac{1}{2v_1^2} (M_{H^0}^2 \cos^2 \alpha + M_{h^0}^2 \sin^2 \alpha), \quad \lambda_2 = \frac{1}{2v_2^2} (M_{H^0}^2 \sin^2 \alpha + M_{h^0}^2 \cos^2 \alpha) \quad (30)$$

$$\lambda_3 = \frac{(M_{H^0}^2 - M_{h^0}^2)}{2v_1 v_2} \sin 2\alpha + 2 \frac{M_{H^\pm}^2}{v^2}, \quad \lambda_4 = \frac{M_{A^0}^2 - 2M_{H^\pm}^2}{v^2} \quad \lambda_5 = -\frac{M_{A^0}^2}{v^2}, \quad (31)$$

$$\lambda_T = (\lambda_3 + \lambda_4 + \lambda_5) = \frac{(M_{H^0}^2 - M_{h^0}^2)}{2v_1 v_2} \sin 2\alpha. \quad (32)$$

Particular cases

To obtain Eq. (30)-(32) we have considered the following relations: Since Eq. (27) is equivalent to

$$M_{H^0, h^0}^2 = \left[\lambda_1 v_1^2 + \lambda_2 v_2^2 \pm (\lambda_1 v_1^2 - \lambda_2 v_2^2) \sqrt{1 + y^2} \right] \quad (33)$$

and

$$\sqrt{1 + y^2} = \sqrt{1 + (\tan 2\alpha)^2} = \frac{1}{\cos 2\alpha}, \quad (34)$$

we obtain from Eq. (25)

$$\sin 2\alpha = \frac{(\lambda_3 + \lambda_4 + \lambda_5) v_1 v_2}{\sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + (\lambda_3 + \lambda_4 + \lambda_5)^2 (v_1 v_2)^2}} \quad (35)$$

as well as

$$\cos 2\alpha = \frac{\lambda_1 v_1^2 - \lambda_2 v_2^2}{\sqrt{(\lambda_1 v_1^2 - \lambda_2 v_2^2)^2 + (\lambda_3 + \lambda_4 + \lambda_5)^2 (v_1 v_2)^2}} \quad (36)$$

and

$$\frac{(M_{H^0}^2 - M_{h^0}^2) \cos 2\alpha}{2} = (\lambda_1 v_1^2 - \lambda_2 v_2^2), \quad \frac{M_{H^0}^2 + M_{h^0}^2}{2} = [\lambda_1 v_1^2 + \lambda_2 v_2^2]. \quad (37)$$

With these equations it is easy to obtain Eq. (30).

a.- In the case when the mixing angle is $\alpha = 0$, i.e., $\lambda_T = 0$, $\lambda_3 > 0$,

$$\lambda_1 = \frac{1}{2v_1^2} M_{H^0}^2, \quad \lambda_2 = \frac{1}{2v_2^2} M_{h^0}^2, \quad \lambda_3 = 2 \left(\frac{M_{H^\pm}}{v} \right)^2, \quad (38)$$

$$\lambda_4 = \left(\frac{M_{A^0}}{v} \right)^2 - 2 \left(\frac{M_{H^\pm}}{v} \right)^2, \quad \lambda_5 = - \left(\frac{M_{A^0}}{v} \right)^2, \quad (39)$$

and

$$M_{H^0}^2 - M_{h^0}^2 = 2(\lambda_1 v_1^2 - \lambda_2 v_2^2). \quad (40)$$

The parameters μ_1, μ_2 in the potential Eq. (3) are related to the neutral Higgs particles in a very simple way, similar to the one between the parameters λ and μ in the SM.

$$-2\mu_1^2 = M_{H^0}^2 = 2\lambda_1 v_1^2, \quad -2\mu_2^2 = M_{h^0}^2 = 2\lambda_2 v_2^2. \quad (41)$$

and the vevs satisfy

$$v^2 = (v_1^2 + v_2^2) = - \left(\frac{\mu_2^2}{\lambda_2} + \frac{\mu_1^2}{\lambda_1} \right) = \frac{1}{2\lambda_1\lambda_2} (\lambda_1 M_{h^0}^2 + \lambda_2 M_{H^0}^2). \quad (42)$$

In this particular case, each Higgs particle is associated with a specific parameter $\lambda_1, \lambda_2, \lambda_3, \lambda_5$.

$$M_{H^0} = v_1 \sqrt{2\lambda_1}, \quad M_{h^0} = v_2 \sqrt{2\lambda_2}, \quad M_{H^\pm} = \frac{v}{\sqrt{2}} \sqrt{\lambda_3}, \quad M_{A^0} = v \sqrt{|\lambda_5|}. \quad (43)$$

The degeneracy in the masses M_{H^0}, M_{h^0} implies that $\lambda_1 v_1^2 - \lambda_2 v_2^2 = 0$.

b.- In the case when the mixing angle is $\alpha = \pi/2$, $\lambda_T = 0$

$$\lambda_1 = \frac{1}{2v_1^2} M_{h^0}^2, \quad \lambda_2 = \frac{1}{2v_2^2} M_{H^0}^2, \quad \lambda_3 = 2 \left(\frac{M_{H^\pm}}{v} \right)^2, \quad (44)$$

$$\lambda_4 = \left(\frac{M_{A^0}}{v} \right)^2 - 2 \left(\frac{M_{H^\pm}}{v} \right)^2, \quad \lambda_5 = - \left(\frac{M_{A^0}}{v} \right)^2. \quad (45)$$

The parameters μ_1, μ_2 in Eq. (12) become:

$$-2\mu_1^2 = M_{h^0}^2 = 2\lambda_1 v_1^2, \quad -2\mu_2^2 = M_{H^0}^2 = 2\lambda_2 v_2^2. \quad (46)$$

and

$$v^2 = \frac{1}{2\lambda_1\lambda_2} (\lambda_1 M_{H^0}^2 + \lambda_2 M_{h^0}^2). \quad (47)$$

As in the former case, each Higgs particle is associated with a specific parameter λ_i , and (H^0, h^0) interchange places.

$$M_{h^0} = v_1 \sqrt{2\lambda_1}, \quad M_{H^0} = v_2 \sqrt{2\lambda_2}, \quad M_{H^\pm} = \frac{v}{\sqrt{2}} \sqrt{\lambda_3}, \quad M_{A^0} = v \sqrt{|\lambda_5|}. \quad (48)$$

In this section, we have considered the main features of various special cases for the parameter α where a decoupling of the Higgs bosons take place, and the case where the masses of the CP-even neutral particles coincide.

V. VACUUM STABILITY CONSTRAINTS

Bounds due to the positive mass-values

Due to the fact that the masses are positive, from the previous results, one gets information for the allowed values of the λ_i parameters.

$$\lambda_1 > 0, \quad \lambda_2 > 0, \quad (\lambda_4 + \lambda_5) < 0, \quad \lambda_5 < 0, \quad \lambda_4 < |\lambda_5|. \quad (49)$$

and Eq. (27) implies that

$$\lambda_1\lambda_2 > \frac{1}{4}(\lambda_3 + \lambda_4 + \lambda_5)^2. \quad (50)$$

In terms of the masses, the conditions in Eq. (49) and Eq. (50) become trivial

$$\lambda_1 = \frac{1}{2v_1^2} [M_{H^0}^2 \cos^2 \alpha + M_{h^0}^2 \sin^2 \alpha] > 0, \quad (51)$$

$$\lambda_2 = \frac{1}{2v_2^2} [M_{H^0}^2 \sin^2 \alpha + M_{h^0}^2 \cos^2 \alpha] > 0, \quad (52)$$

$$\lambda_1\lambda_2 - \frac{1}{4}(\lambda_3 + \lambda_4 + \lambda_5)^2 = \frac{1}{4v_1^2 v_2^2} M_{H^0}^2 M_{h^0}^2 > 0. \quad (53)$$

To improve previous information about the allowed values of the quartic couplings and therefore for the masses, we have explored the consequences of considering the vacuum stability conditions (VSC), through the method of the Lagrange multipliers.

Lagrangian multipliers method and the VSC

Considering one restriction: Let us introduce the variables x_i and the parameters b_i , defined as [13]

$$x_1 = |\Phi_1|^2, \quad x_2 = |\Phi_2|^2, \quad x_3 = \frac{1}{2} (\Phi_1^\dagger \Phi_2 + \Phi_2^\dagger \Phi_1), \quad x_4 = \frac{1}{2i} (\Phi_1^\dagger \Phi_2 - \Phi_2^\dagger \Phi_1), \quad (54)$$

$$b_{11} = \lambda_1, \quad b_{22} = \lambda_2, \quad b_{33} = (\lambda_4 + \lambda_5), \quad b_{44} = (\lambda_4 - \lambda_5), \quad b_{12} = \lambda_3, \quad (55)$$

the potential in Eq. (3) becomes $V = V_0 + x_1^2 F_0$, where $V_0 = \mu_1^2 x_1 + \mu_2^2 x_2$, and

$$F_0 = b_{11} + b_{12}\xi_2 + b_{22}\xi_2^2 + b_{33}\xi_3^2 + b_{44}\xi_4^2, \quad \xi_i = \frac{x_i}{x_1}, \quad i = 2, 3, 4. \quad (56)$$

Using the Cauchy-Schwartz inequality

$$\left| \Phi_1^\dagger \Phi_2 \right|^2 \leq \Phi_1^\dagger \Phi_1 \Phi_2^\dagger \Phi_2 = |\Phi_1|^2 |\Phi_2|^2, \quad (57)$$

we obtain the condition

$$f(\xi_2, \xi_3, \xi_4) = \xi_3^2 + \xi_4^2 - \xi_2 \leq 0. \quad (58)$$

We now introduce the Lagrange multiplier Λ_1 in the quartic sector of the potential [14] related to the condition imposed by the Cauchy-Schwarz inequality Eq. (58)

$$F(\xi_2, \xi_3, \xi_4, \Lambda_1) = F_0(\xi_2, \xi_3, \xi_4) + \Lambda_1 f(\xi_2, \xi_3, \xi_4) = b_{11} + (b_{12} - \Lambda_1)\xi_2 + b_{22}\xi_2^2 + (b_{33} + \Lambda_1)\xi_3^2 + (b_{44} + \Lambda_1)\xi_4^2 \quad (59)$$

and apply the stability (positivity) condition in Eq. (59) after obtaining the derivatives

$$\frac{\partial F}{\partial \xi_2} = \frac{\partial F}{\partial \xi_3} = \frac{\partial F}{\partial \xi_4} = \frac{\partial F}{\partial \Lambda_1} = 0, \quad (60)$$

to evaluate the minimum value for $F(\xi_2, \xi_3, \xi_4, \Lambda_1)$ in the region of interest.

The following equations are to be solved.

$$2b_{22}\xi_2 + b_{12} - \Lambda_1 = 0, \quad 2(b_{33} + \Lambda_1)\xi_3 = 0, \quad 2(b_{44} + \Lambda_1)\xi_4 = 0, \quad \xi_3^2 + \xi_4^2 - \xi_2 = 0. \quad (61)$$

There are two solutions denoted by A_i , $i = 1, 2$, where $\xi_2 = \xi_3^2 + \xi_4^2$, and

$$\Lambda_1 = 2b_{22}\xi_2 + b_{12}, \quad (b_{33} + 2b_{22}\xi_2 + b_{12})\xi_3 = 0, \quad (b_{44} + 2b_{22}\xi_2 + b_{12})\xi_4 = 0, \quad (62)$$

where the stability condition becomes

$$F(\xi_2, \xi_3, \xi_4, \Lambda_1)|_{\min} = (b_{11} - b_{22}\xi_2^2 + (b_{33} + \Lambda_1)\xi_3^2 + (b_{44} + \Lambda_1)\xi_4^2)|_{\min} > 0. \quad (63)$$

There is another solution, case B , in which $\Lambda_1 = 0$, and $\xi_3^2 + \xi_4^2 - \xi_2 \leq 0$. Where

$$F(\xi_2, \xi_3, \xi_4)|_{\min} = (b_{11} + b_{12}\xi_2 + b_{22}\xi_2^2 + b_{33}\xi_3^2 + b_{44}\xi_4^2)|_{\min} > 0. \quad (64)$$

In the first solution A_1 we consider:

$$\xi_4 = 0, \quad \xi_3 \neq 0, \quad \xi_2 = -\frac{1}{2b_{22}}(b_{12} + b_{33}) = \xi_3^2. \quad (65)$$

The conditions and implications for a minimum in the region of interest are:

$$b_{12} + b_{33} \leq 0, \quad b_{22} > 0, \quad \lambda_2 > 0, \quad \lambda_3 + \lambda_4 + \lambda_5 \leq 0. \quad (66)$$

In the second solution A_2

$$\xi_4 \neq 0, \quad (b_{33} + 2b_{22}\xi_2 + b_{12})\xi_3 = 0, \quad b_{44} + 2b_{22}\xi_2 + b_{12} = 0. \quad (67)$$

We have two possibilities, $\xi_3 = 0$, and $\xi_3 \neq 0$, and the existence of a minimum requires, if $\xi_3 = 0$

$$\xi_2 = \xi_4^2, \quad \xi_2 = -\frac{(b_{12} + b_{44})}{2b_{22}}, \quad b_{22} > 0, \quad b_{12} + b_{44} < 0, \quad \lambda_2 > 0, \quad \lambda_3 + \lambda_4 - \lambda_5 \leq 0. \quad (68)$$

If $\xi_3 \neq 0$, the implications are

$$\xi_2 = -\frac{(b_{12} + b_{44})}{2b_{22}} = -\frac{(b_{12} + b_{33})}{2b_{22}}, \quad (b_{33} - b_{44})\xi_3 = 0, \quad \lambda_5 = 0, \quad \lambda_3 + \lambda_4 \leq 0. \quad (69)$$

In case B , the equations to solve are $2b_{22}\xi_2 + b_{12} = 0$, $2b_{33}\xi_3 = 0$, $2b_{44}\xi_4 = 0$, and the conditions to have a minimum with its implications are

$$\xi_2 = -\frac{b_{12}}{2b_{22}} > 0, \quad b_{12} \leq 0, \quad b_{22} > 0, \quad \lambda_3 < 0. \quad (70)$$

Now we apply the stability condition in Eq. (63) and obtain in cases A_1 and A_2

$$4b_{22}b_{11} \geq (b_{12} + b_{33})^2 \Rightarrow -2\sqrt{\lambda_1\lambda_2} < (\lambda_3 + \lambda_4 + \lambda_5), \quad (71)$$

$$4b_{11}b_{22} > (b_{12} + b_{44})^2 \Rightarrow -2\sqrt{\lambda_1\lambda_2} < (\lambda_3 + \lambda_4 - \lambda_5), \quad (72)$$

$$b_{33} = b_{44}, \quad \lambda_5 = 0 \Rightarrow -2\sqrt{\lambda_1\lambda_2} < (\lambda_3 + \lambda_4). \quad (73)$$

Now, in case B and Eq. (64), the result is

$$4b_{22}b_{11} > (b_{12})^2 \Rightarrow \lambda_3 > -2\sqrt{\lambda_1\lambda_2}. \quad (74)$$

Performing the second derivative

$$F(\xi_2, \xi_3, \xi_4, \Lambda_1) = b_{11} + (b_{12} - \Lambda_1)\xi_2 + b_{22}\xi_2^2 + (b_{33} + \Lambda_1)\xi_3^2 + (b_{44} + \Lambda_1)\xi_4^2. \quad (75)$$

We obtain

$$\frac{\partial}{\partial \xi_2} \left(\frac{\partial F}{\partial \xi_2} \right) = \frac{\partial}{\partial \xi_2} (2b_{22}\xi_2 + b_{12} - \Lambda_1) = 2b_{22} > 0. \quad (76)$$

After imposing the stability condition, with one restriction we obtain these new boundary values for the quartic couplings:

$$-2\sqrt{\lambda_1\lambda_2} < (\lambda_3 + \lambda_4 + \lambda_5), \quad -2\sqrt{\lambda_2\lambda_1} \leq \lambda_3. \quad (77)$$

Considering now two restrictions: Following the same method as in the former case, considering now the conditions

$$x_3^2 + x_4^2 \leq x_1x_2, \quad x_1 + x_2 - v^2 = 0 \quad (78)$$

and two Lagrange multipliers, the function in consideration is:

$$F_{2c}(x_1, x_2, x_3, x_4, \Lambda_1, \Lambda_2) = b_{11}x_1^2 + b_{22}x_2^2 + b_{33}x_3^2 + b_{44}x_4^2 + b_{12}x_1x_2 + \Lambda_1(x_3^2 + x_4^2 - x_1x_2) + \Lambda_2(x_1 + x_2 - v^2) \quad (79)$$

The new equations to solve are

$$\begin{aligned} 2b_{11}x_1 + b_{12}x_2 + \Lambda_2 - \Lambda_1x_2 &= 0, & 2b_{22}x_2 + b_{12}x_1 + \Lambda_2 - \Lambda_1x_1 &= 0, \\ (b_{33} + \Lambda_1)x_3 &= 0, & (b_{44} + \Lambda_1)x_4 &= 0. \\ \Lambda_1(x_3^2 + x_4^2 - x_1x_2) &= 0, & \Lambda_2(x_1 + x_2 - v^2) &= 0. \end{aligned} \quad (80)$$

For $F_{2c}(x_1, x_2, x_3, x_4, \Lambda_1, \Lambda_2)|_{\min}$

$$x_1 = \left(\frac{b_{12} - \Lambda_1 - 2b_{22}}{4b_{22}b_{11} - (b_{12} - \Lambda_1)^2} \right) \Lambda_2 > 0, \quad x_2 = \left(\frac{b_{12} - \Lambda_1 - 2b_{11}}{4b_{22}b_{11} - (b_{12} - \Lambda_1)^2} \right) \Lambda_2 > 0. \quad (81)$$

and

$$F_{2c}(x_1, x_2, x_3, x_4, \Lambda_1, \Lambda_2)|_{\min} = b_{11}x_1^2 + b_{22}x_2^2 + (b_{12} - \Lambda_1)x_1x_2|_{\min} \quad (82)$$

The solutions satisfy the following requirements:

$$\frac{F_{2c}(x_1, x_2, x_3, x_4, \Lambda_1, \Lambda_2)|_{\min}}{\Lambda_2^2} = \frac{(b_{11} + b_{22} - b_{12} + \Lambda_1)}{4b_{22}b_{11} - (b_{12} - \Lambda_1)^2} > 0, \quad (83)$$

with

$$\Lambda_2 = \frac{4b_{22}b_{11} - (b_{12} - \Lambda_1)^2}{2(b_{12} - \Lambda_1 - b_{22} - b_{11})}v^2 < 0. \quad (84)$$

The requirements in Eqs. (81) imply

$$\begin{aligned} (b_{12} - \Lambda_1)^2 - 4b_{11}b_{22} &< 0, & b_{12} - \Lambda_1 - 2b_{22} &< 0, \\ (b_{12} - \Lambda_1 - b_{11} - b_{22}) &< 0, & b_{12} - \Lambda_1 - 2b_{11} &< 0. \end{aligned} \quad (85)$$

The additional solutions that come from Eqs. (80) yield

$$\Lambda_1 = -b_{33}, \quad \Lambda_1 = -b_{44}, \quad b_{33} = b_{44}.$$

Then

$$(\lambda_3 + \lambda_4 \pm |\lambda_5|)^2 < 4\lambda_1\lambda_2, \quad \lambda_3 + \lambda_4 \pm |\lambda_5| < \lambda_1 + \lambda_2, \quad \lambda_3 + \lambda_4 \pm |\lambda_5| < 2\lambda_2, \quad \lambda_3 + \lambda_4 \pm |\lambda_5| < 2\lambda_1$$

and

$$-2\sqrt{\lambda_1\lambda_2} < (\lambda_3 + \lambda_4 \pm |\lambda_5|), \quad \lambda_3 + \lambda_4 \pm |\lambda_5| < 2\lambda_1, \quad \lambda_3 + \lambda_4 \pm |\lambda_5| < 2\lambda_2.$$

Thus we obtain previous results plus

$$\lambda_1 + \lambda_2 > \lambda_3 + \lambda_4 \pm |\lambda_5|, \quad \lambda_1 + \lambda_2 > \lambda_3 + \lambda_4. \quad (86)$$

Bounds from extreme stability conditions

Let us now determine the behavior of the quartic couplings in the case where the quartic Higgs potential V_4 has its lowest possible value. In correspondence with Eq. (3) using the notation of Eq. (54) when $x_1 = v_1^2$, $x_2 = v_2^2$ the V_4 can be simplified as follows

$$V_4 = \lambda_1x_1^2 + \lambda_2x_2^2 + \lambda_Tx_1x_2 = \left(\sqrt{\lambda_1}x_1 - \sqrt{\lambda_2}x_2 \right)^2 + \left(\lambda_T + 2\sqrt{\lambda_1\lambda_2} \right)x_1x_2 \geq 0.$$

In the Extreme case, the condition to be satisfied is $V_4 = 0$, then

$$\sqrt{\frac{\lambda_1}{\lambda_2}} = \frac{x_2}{x_1} = \frac{v_2^2}{v_1^2}, \quad \lambda_T = -2\sqrt{\lambda_1\lambda_2}, \quad \Rightarrow \quad \frac{\lambda_1}{\lambda_2} = \frac{v_2^4}{v_1^4} = (\tan\beta)^4.$$

In this case the Higgs masses become

$$M_{H^0} = \begin{cases} (4\lambda_1\lambda_2)^{1/4} v, & \lambda_1 \neq \lambda_2, \quad v_1 \neq v_2, \\ (2\lambda)^{1/2} v, & \lambda_1 = \lambda_2 = \lambda, \quad v_1 = v_2 = v/\sqrt{2}. \end{cases}$$

In another interesting case, which is the Semi-extreme case, the $V_4 = (\sqrt{\lambda_1}x_1 - \sqrt{\lambda_1}x_1)^2 > 0$, and

$$\frac{\lambda_1}{\lambda_2} \neq (\tan \beta)^4, \quad \lambda_T = -2\sqrt{\lambda_1\lambda_2},$$

the M_{H^0} becomes

$$M_{H^0} = \sqrt{2} (\lambda_1 v^2 + (\lambda_2 - \lambda_1) v_2^2)^{1/2}.$$

In both cases

$$M_{h^0} = 0, \quad M_{H^\pm} = \left(\frac{1}{2} |\lambda_4 + \lambda_5| \right)^{1/2} v, \quad M_{A^0} = |\lambda_5|^{1/2} v.$$

Numerical evaluation of the Higgs masses in terms of $\tan \beta$

In general, according to Eqs. (29)-(31) the mass dependence on v_1, v_2 , or $v = \sqrt{v_1^2 + v_2^2}$, can be reformulated in terms of v and v_2 . Therefore, with known v one can plot those masses in terms of v_2 , and determine their dependence on v_2 or $\tan \beta$ ($\tan \beta = v_2/\sqrt{v^2 - v_2^2}$), as in Fig. 1.

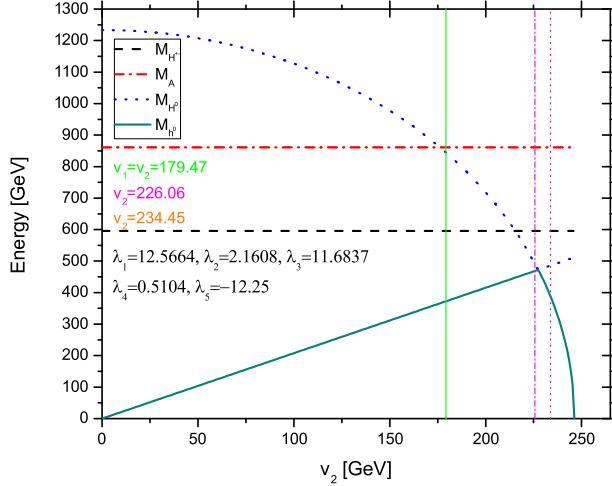


FIG. 1. The v_2 dependence of the Higgs masses for various sets of quartic couplings.

As the favoured Standard Model Higgs mass window is still open, and as the 2HDM Higgses are nonstandard, almost none of the masses range has yet been ruled out, one can explore several hypothetical situations and analyze the consequences of some published values for M_{H^\pm} .

We will now proceed to numerically evaluate the Higgs masses under different conditions for the quartic parameters at the energy scale $E = M_t$, where M_t is the mass of the quark top. First we will reproduce the M_{H^\pm} value given in [15], considering several sets of λ 's. We choose $\lambda_4 = -5.52$, and $\lambda_5 = -6.0$ from many combinations that give $M_{H^\pm} = 609.1, M_{A^0} = 621.7$. As we shall see, in the extreme case, with one of the symmetries considered in [16] in which $\lambda_1 = \lambda_2 = \lambda$, and $\tan \beta = 1$, or $\tan \beta = 5$ all the Higgs masses acquire constant values and $M_{h^0} = 0$. With the previous λ_4, λ_5 , we obtain for H^0 values which are not ruled out experimentally (for the SM Higgs) according to [17], in the following way: ($\lambda = 0.485, \lambda_3 = 10.55, \lambda_T = -0.97, M_{H^0} = 250$), ($\lambda = 0.653, \lambda_3 = 10.21, \lambda_T = -1.30, M_{H^0} = 290$) and ($\lambda = 1.94, \lambda_3 = 7.64, \lambda_T = -3.88, M_{H^0} = 600$).

Though the masses do not depend on $\tan \beta$ at this energy scale, we will explore, in the next section, their behavior and dependence on it at higher energies scales.

Let us now classify the several cases, we will consider , in terms of the different stability conditions for the λ_i s
A.- Extreme case in which both equalities are satisfied

$$\lambda_T = -2\sqrt{\lambda_1\lambda_2}, \quad \cup \quad \lambda_1/\lambda_2 = (\tan \beta)^4 \quad (87)$$

B1.- Semi-extreme case:

$$\lambda_T = -2\sqrt{\lambda_1\lambda_2}, \quad \lambda_1/\lambda_2 \neq (\tan \beta)^4 \quad (88)$$

B2.- Semi-extreme case:

$$\lambda_T \neq -2\sqrt{\lambda_1\lambda_2}, \quad \lambda_1/\lambda_2 = (\tan \beta)^4 \quad (89)$$

C.- Lagrange inequality condition

$$\lambda_T \geq -2\sqrt{\lambda_1\lambda_2} \quad (90)$$

D.- Yukawa - Universality condition

$$\tan \beta = M_t/M_b, \quad g_t = g_b \quad (91)$$

The former cases will be combined with two additional conditions for the quartic couplings:

Case 1 : $\lambda_1 = \lambda_2$, and Case 2: $\lambda_1 \neq \lambda_2$.

In case 1A , 2A and 1B, all the masses have a constant value in the interval $0 \leq v_2 \leq v$ due to their only dependence on v , but in spite of the an explicit independence of the masses on v_2 , $\tan \beta$ plays an important role on the energy scale dependence of the masses, as we shall see in the numerical solutions of the Renormalization Group Equations.which is fixed. In case 2B, M_{H^0} depends explicitly on v_2 and therefore on $\tan \beta$.

To compare with the masses $M_{H^\pm} = 609$, $M_A = 621.7$ given in Ref.[15], we consider three kinds of compositions of the λ_i $i = 1, \dots, 5$ parameters as in a), b) and c), which produce seven cases. The properties of these cases are to be analized in the following section.

a.-When $\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ 0.485 & 0.485 & 10.55 & -5.52 & -6.0 \end{pmatrix}$, which means that $\lambda_T = -2\sqrt{\lambda_1\lambda_2} = -0.97$ and $(\lambda_1/\lambda_2)^{1/4} = 1.0$, we obtain

Case	$\tan \beta$	M_{h^0}	M_{H^0}	v_2
1A	1.0	0	250	179.47
1D	41.2	0	250	253.7

b.- When $\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ 6.0 & 1/6 & 9.51 & -5.51 & -6.0 \end{pmatrix}$, i.e., $\lambda_T = -2\sqrt{\lambda_1\lambda_2} = -2.0$ and $(\lambda_1/\lambda_2)^{1/4} = 2.45$, we obtain

Case	$\tan \beta$	M_{h^0}	M_{H^0}	v_2
2A	2.45	0	358.9	234.9

c.- When $\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ 12.0 & 3.0 & 11.5 & -5.5 & -6.0 \end{pmatrix}$ with $\lambda_T = 0.0$, $-2\sqrt{\lambda_1\lambda_2} = -12.0$ and $(\lambda_1/\lambda_2)^{1/4} = 1.4$, we explore different $\tan \beta$ cases

Case	$\tan \beta$	M_{h^0}	M_{H^0}	v_2
2B2	1.4	507.6	717.9	207.2
2C	2.0	556.0	556.0	227.0
2C	5.0	243.9	609.6	248.9
2D	41.2	30.2	621.5	253.7

d.- Considering now smaller values for $\{M_{H^\pm}, M_A\} = \{253.8, 240.8\}$, with interesting properties for the energy range of validity of the 2HDM, with $\lambda_T = -2\sqrt{\lambda_1\lambda_2} = -0.97$ and $(\lambda_1/\lambda_2)^{1/4} = 1.0$, for $\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ 0.48 & 0.48 & 1.03 & -1.1 & -0.9 \end{pmatrix}$

Case	$\tan \beta$	M_{h^0}	M_{H^0}	v_2
1A	1.0	0	250	179.5
1D	41.2	0	250	253.7

e.- Finally, for even lower masses $\{M_{H^\pm}, M_A\} = \{170.3, 160.5\}$, which arise from $\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ 0.125 & 0.125 & 0.65 & -0.5 & -0.4 \end{pmatrix}$, where $\lambda_T = -2\sqrt{\lambda_1\lambda_2} = -0.25$ and $(\lambda_1/\lambda_2)^{1/4} = 1.0$, we obtain

Case	$\tan\beta$	M_{h^0}	M_{H^0}	v_2
1D, 1B1	41.2	0	127	253.7

VI. TRIVIALITY CONSTRAINTS.

Renormalization group equations

In this section we explore the asymptotic behavior of the parameters in the model, and their relations, through the Renormalization Group Equations (RGE) [8]. The RGE are a powerful tool to determine by the triviality principle, the energy bounds of the parameters and the validity of the model. In order to proceed in this way, to numerically evaluate the energy dependence of the λ_i quartic couplings, it is necessary to consider the RGE of all the parameters, i.e., the gauge group couplings g_1, g_2, g_3 of the symmetry groups $U(1), SU(2), SU(3)$, the vacuum expectation values v_1, v_2 , and the Yukawa couplings of the top and the down quark sectors g_t and g_b respectively Refs. [18].

The RGE determine the dependence of the coupling constants and other parameters of the Lagrangian on t , defined as $t = \ln(E/m_t)$, where E is the renormalization point energy. The RGE for the gauge couplings g_1, g_2, g_3 are:

$$\frac{dg_k}{dt} = \frac{1}{(4\pi)^2} b_k g_k^3 \quad (i = 1, 2, 3), \quad (92)$$

The RGE for the Yukawa couplings of the top and bottom quarks g_t, g_b are

$$\frac{dg_t}{dt} = \frac{1}{(4\pi)^2} \left(\frac{9}{2} g_t^2 + \frac{1}{2} g_b^2 - \left(\frac{17}{20} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) g_t \right) g_t, \quad (93)$$

$$\frac{dg_b}{dt} = \frac{1}{(4\pi)^2} \left(\frac{9}{2} g_b^2 + \frac{1}{2} g_t^2 - \left(\frac{1}{4} g_1^2 + \frac{9}{4} g_2^2 + 8g_3^2 \right) g_b \right) g_b \quad (94)$$

and for the vacuum expectation values v_1 and v_2

$$\frac{d}{dt} v_1 = \frac{1}{(4\pi)^2} \left[-3g_t^2 + ((9/20)g_1^2 + (9/4)g_2^2) \right] v_1, \quad (95)$$

$$\frac{d}{dt} v_2 = \frac{1}{(4\pi)^2} \left[-3g_b^2 + ((9/20)g_1^2 + (9/4)g_2^2) \right] v_2 \quad (96)$$

In the equations for the quartic couplings we include the quark Yukawa contributions of both sectors.

$$\begin{aligned} \frac{d\lambda_1}{dt} &= \frac{1}{16\pi^2} \left\{ 24(\lambda_1)^2 - 3\lambda_1 \left[3g^2 + (g')^2 - 4g_t^2 \right] + 2(\lambda_3)^2 + (\lambda_4)^2 + (\lambda_5)^2 + 2\lambda_3\lambda_4 + \frac{3}{8} \left(g^2 + (g')^2 \right)^2 + \frac{3}{4}g^4 - 6g_t^4 \right\}, \\ \frac{d\lambda_2}{dt} &= \frac{1}{16\pi^2} \left\{ 24(\lambda_2)^2 - 3\lambda_2 \left[3g^2 + (g')^2 - 4g_b^2 \right] + 2(\lambda_3)^2 + (\lambda_4)^2 + (\lambda_5)^2 + 2\lambda_3\lambda_4 + \frac{3}{8} \left[(g')^2 + g^2 \right]^2 + \frac{3}{4}g^4 - 6g_b^4 \right\}, \\ \frac{d\lambda_3}{dt} &= \frac{1}{16\pi^2} \left\{ 4(\lambda_3)^2 + 4(3\lambda_3 + \lambda_4)(\lambda_1 + \lambda_2) - 3\lambda_3 \left[3g^2 + (g')^2 - 2(g_t^2 + g_b^2) \right] \right. \\ &\quad \left. + 2(\lambda_4)^2 + 2(\lambda_5)^2 + \frac{3}{4} \left[g^2 - (g')^2 \right]^2 + \frac{3}{2}g^4 - 12g_t^2g_b^2 \right\}, \end{aligned} \quad (97)$$

$$\frac{d\lambda_4}{dt} = \frac{1}{16\pi^2} \left\{ 4(\lambda_4)^2 + 4\lambda_4(\lambda_1 + \lambda_2 + 2\lambda_3) - 3\lambda_4 \left[3g^2 + (g')^2 - 2(g_t^2 + g_b^2) \right] + 8(\lambda_5)^2 + 3g^2(g')^2 + 12g_t^2g_b^2 \right\},$$

$$\frac{d\lambda_5}{dt} = \frac{1}{16\pi^2} \lambda_5 \left\{ 4(\lambda_1 + \lambda_2 + 2\lambda_3 + 3\lambda_4) - 3 \left[3g^2 + (g')^2 - 2(g_t^2 + g_b^2) \right] \right\}.$$

The former equations are the coupled, non linear, ordinary differential equations whose solution provides the information about the renormalization point energy dependence of the masses of the five Higgs particles of the 2HDM. To numerically solve the RGE, the initial or final conditions for the parameters have to be previously chosen. In order to do so we use Ref. [2]. The range of values, we take, for the energy and the variable t are $(E_0 = M_t, E_u) = (173.2, 1.234 \cdot 10^{13})$,

$(t_0 = 0, t_u = 25)$ respectively, where M_t stands for the mass of the quark top and E_u corresponds to the electroweak unification energy where $g_1(E_t) = g_2(E_t)$. The gauge couplings $(g_1, g_2, g_3)_{E=M_t} \simeq (0.4627, 0.6466, 1.2367,)$ are obtained using the following relations

$$\begin{aligned}\alpha_e(M_t) &= g_e^2/4\pi = 1/(127.9), \quad g_1(M_t) = \sqrt{5/3}g_e/\cos\theta_W \\ g_2(M_t) &= g_e/\sin\theta_W, \quad g_3(M_t) = \sqrt{4\pi\alpha_s(M_t)}\end{aligned}$$

where θ_W is the Weinberg angle and $\sin^2\theta_W(M_t) = 0.235$ and $\alpha_s = 0.1217$. The vev standard value that arises from

$$v = 2M_z/\sqrt{g_2^2 + g_e^2}$$

is $v(M_t) = 253.81$ GeV at $M_z = 91.19$ GeV.

In order to specify more rigorously the energy limits for the quartic couplings, we have numerically solved the RGE for the gauge group couplings g_1, g_2, g_3 , (Fig. 2), the vacuum expectation values v_1, v_2 , and the top and the down quark Yukawa couplings g_t and g_d , under the following assumptions:

- The heaviest quark masses are related with the vevs v_1 and v_2 and the Yukawa couplings $g^{(u)}$ and $g^{(d)}$

$$M_t = \frac{v_2}{\sqrt{2}}g_t, \quad \tan\beta = \frac{v_2}{v_1}, \quad M_b = \frac{v_1}{\sqrt{2}}g_b. \quad (98)$$

- The gauge bosons masses are related with the gauge couplings g' and g

$$M_W = \frac{1}{2}vg, \quad M_Z = \frac{M_W}{\cos\theta_W} = \frac{1}{2}v\sqrt{g^2 + (g')^2}, \quad (99)$$

where θ_W is the Weinberg angle and e the electron charge

$$e = g\sin\theta_W = g'\cos\theta_W. \quad (100)$$

- Unification of the Yukawa couplings at $E = M_t$ or at E_u , i.e., $g_b = g_t$, and $\tan\beta = M_t/M_b$.

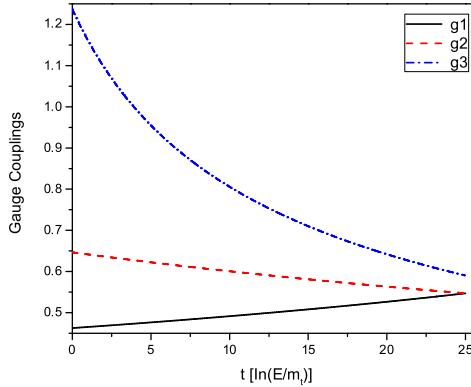


FIG. 2. The energy dependence of the gauge couplings in the 2HDM.

It is interesting to explore now, the energy bounds of the 2DHM, through the running of the quartic couplings which determine the mass values of the Higgses. In the case (c) considered in the previous section, when $M_{H^\pm} = 609$, $M_A = 621.7$, the range of validity of the model is very short $M_t < E < 292$ i.e., $0 < t < 0.52$ as can be seen in Figs. 3–6. There is an intermediate class of the models depicted at Figs. 7, 8, which have an intermediate range of validity $0 < t \lesssim 11$. So we will rather focus our attention on the cases where we can explore the universality of the Yukawa couplings and its unification, to study the mass-hierarchy problem. In this case, as can be seen in Figs. 9–11, the 2HDM is valid in the whole range of energies, this means $M_t < E < E_u$ where E_u is the electroweak unification energy. In Fig. 9 we observe very slow dependence of the quartic couplings and the Higgs masses on the renormalization point energy. This model is characterized by rather small values of the quartic Yukawa couplings and the value of $\tan\beta$ such that it permits the unification of the Yukawa couplings of the up and down quarks $g_t = g_b$. In Figs. 10, 11 we show the evolution of the Yukawa couplings, quartic couplings and the Higgs masses for the case when the Yukawa couplings are unified. In Fig. 10 we assume that they are unified at low energy and in Fig. 11 they are unified at high energy. The evolution of the quartic couplings and Higgs masses are similar in both cases.

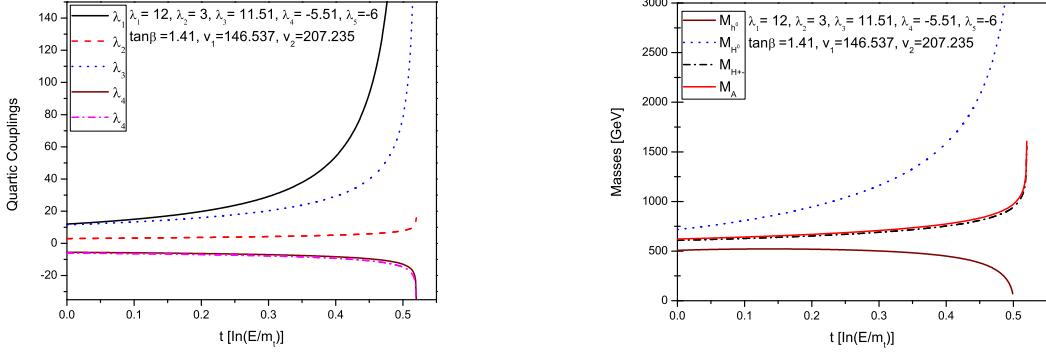


FIG. 3. The energy dependence of the quartic couplings and Higgs masses, case 2B2, with $\tan \beta = 1.41$.

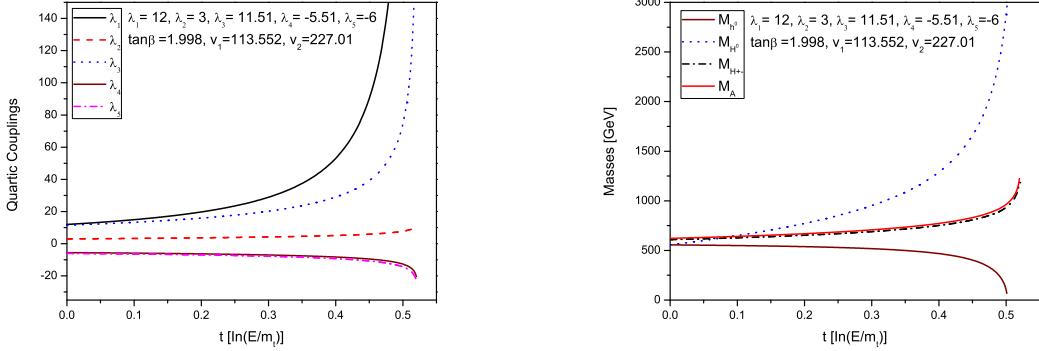


FIG. 4. The energy dependence of the quartic couplings and the Higgs masses, case 2C with $\tan \beta = 2.0$.

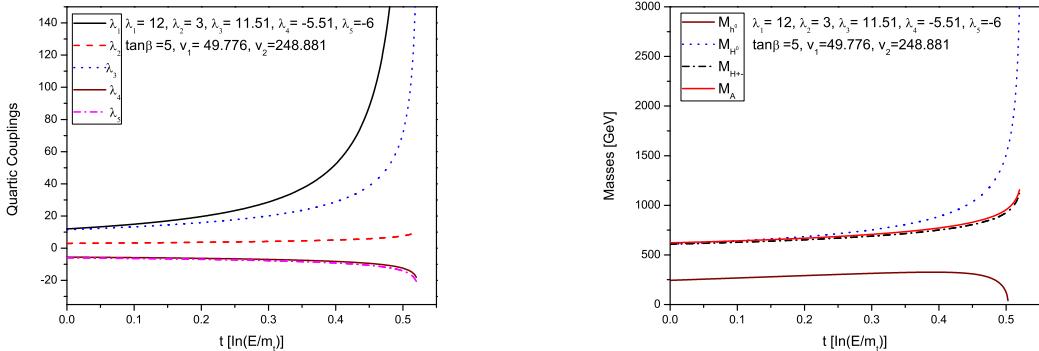


FIG. 5. The energy dependence of the quartic couplings and the Higgs masses, case 2C with $\tan \beta = 5.0$.

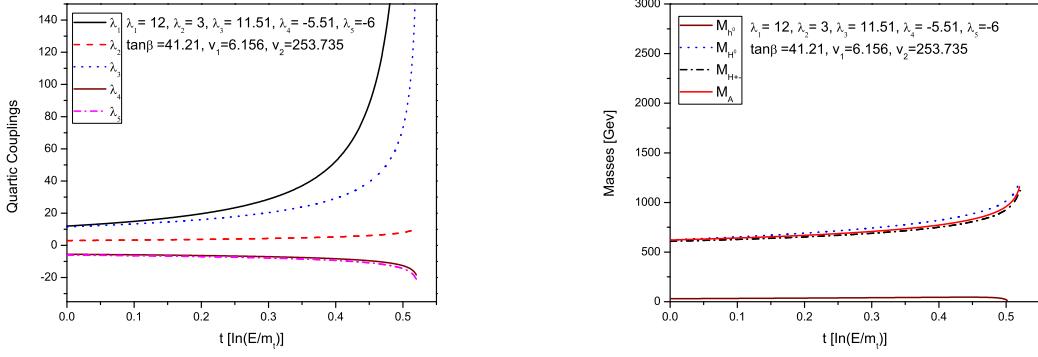


FIG. 6. The energy dependence of the quartic couplings and the Higgs masses, case 2D with $\tan\beta = 41.2$.

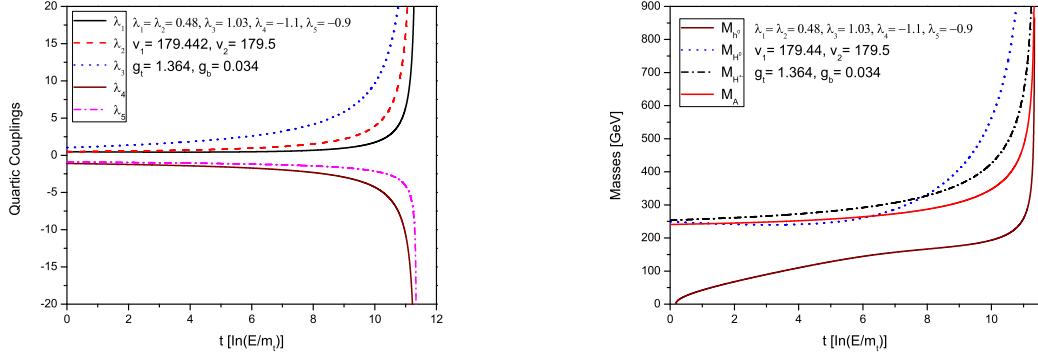


FIG. 7. The energy dependence of the quartic couplings and the Higgs masses, case 1A with $\tan\beta = 1$ (2D) case.

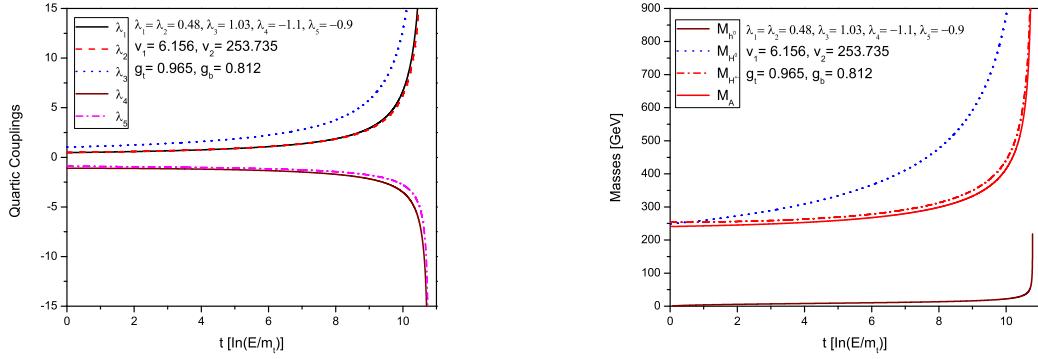


FIG. 8. The energy dependence of the quartic couplings and the Higgs masses, case 1D with $\tan\beta = 41.2$ (2D) case.

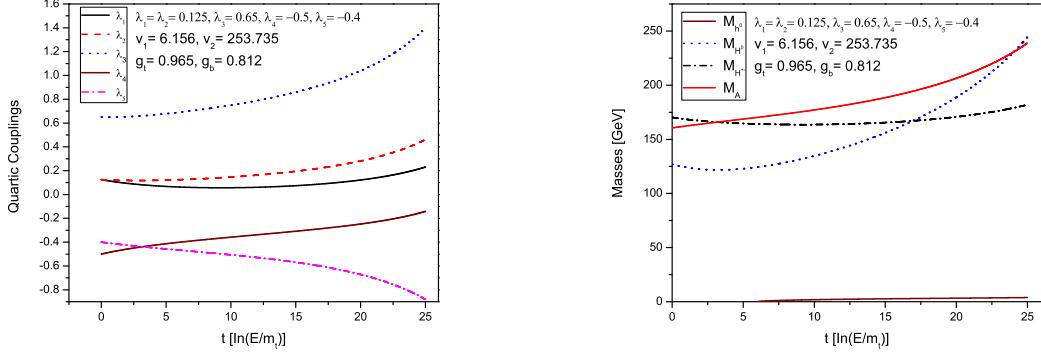


FIG. 9. The energy dependence of the quartic couplings and the Higgs masses, case 1B1 with $\tan \beta = 41.2$.

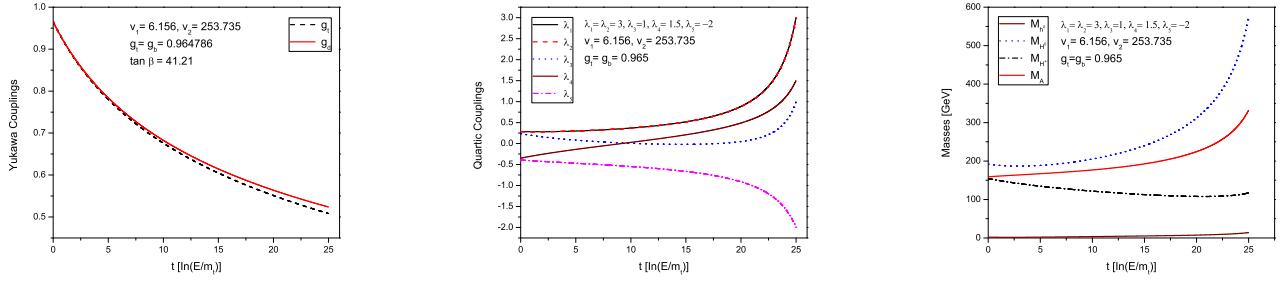


FIG. 10. The energy dependence of the Yukawa couplings, quartic couplings and the Higgs masses in the $\tan \beta = 41.2$ (1D) case with Yukawa couplings $g_t = g_b$ at low energy.

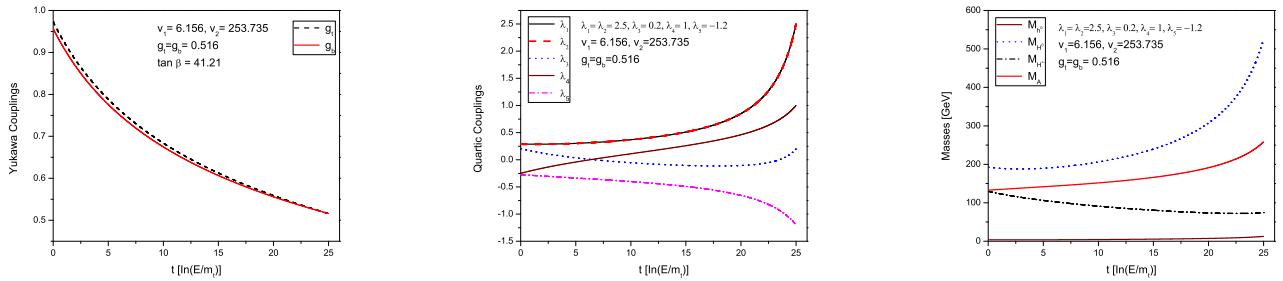


FIG. 11. The energy dependence of the Yukawa couplings, quartic couplings and the Higgs masses in the $\tan \beta = 41.2$ (1D) case with equal Yukawa couplings $g_t = g_b$ at high energy.

VII. RESULTS AND CONCLUSIONS

With the aim to explore the Higgs mass content of the 2HDM extension of the standard model, among the different forms of the Lagrangian describing the same physical reality, we have chosen a specific one, in which the vacuum expectation values of both Higgs fields are real, and for simplicity also preserving the CP symmetry. We have deduced, in this model, the analytical expressions for the masses of the five predicted physical Higgs particles, and expressed the Higgs potential in terms of those masses, using Eqs. (30) and (31). We have also obtained, through the mass formulas, a set of constraints to be satisfied by the scalar parameters that determine the couplings and self-couplings of the Higgs fields introduced in the potential Eq. (3), and through the vacuum stability principle plus the Lagrange Multipliers method, and obtained additional conditions to be satisfied by those couplings.

$$\lambda_1 > 0, \lambda_2 > 0, \quad 4\lambda_1\lambda_2 > (\lambda_3 + \lambda_4 + \lambda_5)^2, \quad (\lambda_4 + \lambda_5) < 0, \quad \lambda_5 < 0, \quad \lambda_4 < |\lambda_5|, \quad (101)$$

and

$$\lambda_1 + \lambda_2 > \lambda_3 + \lambda_4 + \lambda_5, \quad \lambda_3 + \lambda_4 + \lambda_5 + 2\sqrt{\lambda_1\lambda_2} > 0, \quad \lambda_3 + 2\sqrt{\lambda_1\lambda_2} > 0, \quad 4(\lambda_1\lambda_2) \neq \lambda_3^2. \quad (102)$$

We have also looked upon extreme and semiextreme conditions on the Higgs potential and gave a classification of the different cases we analyzed under the RGE.

As many authors base their calculations in symmetry conditions, such as $\lambda_1 = \lambda_2$ and others in a phenomenological study of special events, it is important to analyze the consequences of such assumptions and we tried at least partially address this problem.

There is a batch of data to be analyzed right now in search of some of the favored mass region, and all of it should be examined in the near future. The results of this paper may shed some light on physics of the Higgs sector depending on the properties of the Higgs particle.

We have considered here, symmetries in the λ_i parameters, universality of the Yukawa couplings at low energy E_0 (M_t scale) or high energy E_u (weak-unification scale), hierarchy of the quark masses and the energy range of validity of the model. The other symmetry considered here is the unification of the Yukawa couplings. It seems this symmetry makes the Higgs sector very stable as can be seen in Fig. 9.

In summary, the results in this paper may be a basis for further investigation in relation to the behavior and energy dependent characteristics of the Higgs particles.

We finish by allegorically saying, that our paper still contains a “blank page”, which can only be filled after the discovery of Higgs bosons.

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